



Characterising W radiation in JET-ILW plasmas

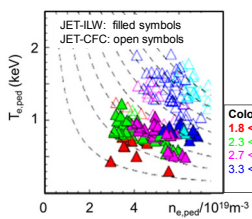
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MOTIVATION FOR STUDY: pedestal in JET-ILW colder tan in JET-C

- JET-ILW (with W divertor): typically $T_{e,ped} \sim 0.7-1.2$ keV. JET-C (CFC divertor) $T_{e,ped} \sim 1.5-2$ keV [1,2]
- This is due in part to the much larger fuelling rates required to maintain the ELM frequency up and avoid W events and accumulation. At larger fuelling, pedestals have larger n_e , lower T_e



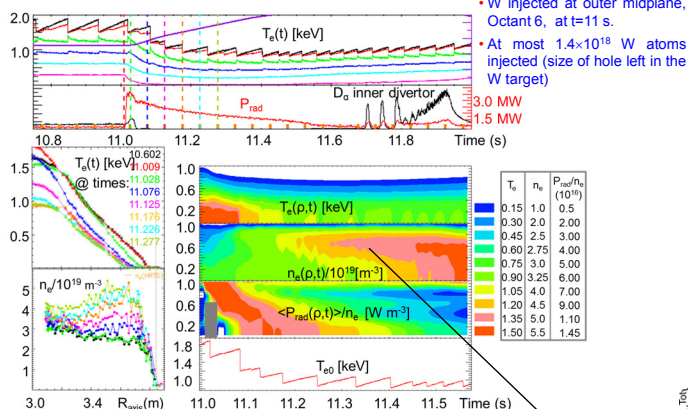
- That explanation is insufficient: at same fuelling levels, pedestals remain colder, particularly at higher densities.
- Neoclassical transport: H-mode pedestal can be W accumulator. W is driven in by inward pinch from ∇n_e , but once it reaches the top of ∇n_e , it would penetrate much more slowly, by diffusion.

Colours:
 1.8 < I_p (MA) < 2.3
 2.3 < I_p (MA) < 2.7
 2.7 < I_p (MA) < 3.3
 3.3 < I_p (MA) < 3.8
 I_p (MA) > 3.8

W can accumulate, in between ELMs, near $T_{e,ped}$. W-associated radiation would slow down the rise of $T_{e,ped}$ in between ELMs, while density continues to rise until an ELM is triggered, resulting in colder pedestals.

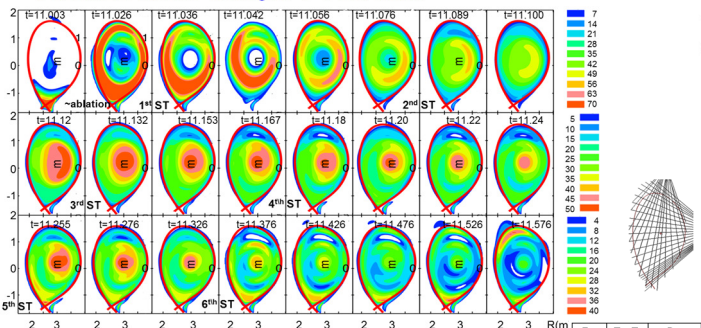
W Ablation into L-mode plasma: JET 90472, 2 MA, 2.4 T, 1.2 MW NBI

To investigate to what extent the difference in $T_{e,ped}$ is due to W radiation in the pedestal region we injected W (by laser ablation) into cold L-mode plasmas at JET.



- T_e collapses at plasma edge: radiating mantle.
- Below 100 eV radiation dominated by D, Be.
- Power detachment, collapse of $T_{e,edge} \rightarrow T_{e,core}$ drop
- Langmuir probes: plasma detaches as soon as W reaches plasma edge, $T_{e,SOL}$ drops from 13 to 5 eV.
- Hollow n_e profile would prevent penetration of W into $p < 0.6$.
- D, V adjusted to simulate effect of sawteeth bringing W inward and force $n_{W,p}(p,t)$ at 11.2 s.

BOLOMETRY, time averaged +/- 5 ms, kW/m²



Persistent poloidal structures and plumes

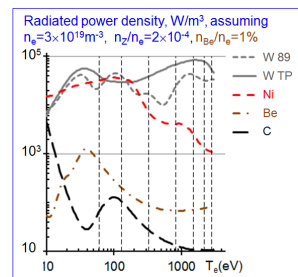
Centrifugal effects produce in-out asymmetry: $n=1$ MHD, 500 Hz: Mach ~ 0.3
 W ions propagate along field line by collisional diffusion, PS regime.
 Random walk estimate of time required for W to spread poloidally
 $t = L^2/D_{||}$
 With $L = \pi(qR+a)$, $q=3$, $n_e = 3 \times 10^{19} \text{ m}^{-3}$ typical times are $\sim 7-85$ ms
 These are comparable to radial transport times.

Z _W	T _e =T _i eV	D 10 ⁴ m ² /s	time ms
W ¹⁸	160	1.1	85
W ^{19a}	210	1.4	67
W ^{19b}	290	1.9	51
W ^{19c}	420	2.8	34
W ^{19d}	500	3.4	28
W ^{20a}	685	5.1	18
W ^{20b}	881	7.0	14
W ^{20c}	1230	13.0	7

W RADIATION, COOLING FUNCTIONS L_Z(T_e)

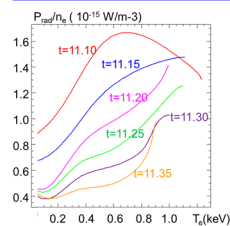
$$P_{rad} = \sum n_e n_Z L_Z(T_e)$$

W is generally a very good radiator in the 0.5-2 keV range. Calculation is challenging. One way of representing the information is via the cooling functions $L_Z(T_e)$.



For low and medium Z impurities (C, Be, Ni) radiation is a decreasing function of T_e from 250 eV: burn through.

For W, two ADAS-based curves are shown:
 W 89: ADAS 89 baseline data, with only low level configurations used to estimate line radiated power;
 W TP: similar to W 89, but with a different choice of configurations, adjusted recombination coefficients to match ASDEX measurements in W²⁴-W⁶⁺ range and more sophisticated ionization rates [3,4].
 Both curves have regions of positive slope: when plasma is heated radiation increases, and it is harder for T_e to rise. But both $L_{W,89}(T_e)$ and $L_{W,TP}(T_e)$ omit important effects, particularly below 1 keV.

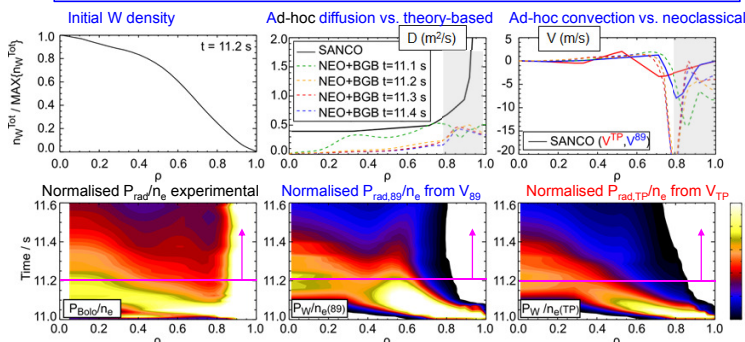


Experimental observations indicate major flaw in dielectronic recombination (DR) coefficients for mid-range ions with an open 4f shell in the ground state. This can increase DR rates in W by a factor of 10. It would significantly move the ionisation balance to higher T_e with the expectation that structure would emerge in $L_{W,TP}(T_e)$ near ~ 600 eV. Work is in progress to incorporate these effects

As a first cut on the data, we plot P_{rad}/n_e as a function of T_e from 11.1 s. The n_W profile is convolved with these curves, but excepting the first timeslice, all exhibit an increasing character from $T_e > 0.6$ keV, steeper above 0.8 keV, and a bump at 0.4 keV, possibly revealing some structure in $L_W(T_e)$.

TRANSPORT MODELLING: radiation from different W models

Ad-hoc D and V used in SANCO to compute P_{rad} from experimental n_e, T_e and two sets of atomic data, like W 89 and W TP.
 Modelling starts at $t=11.2$, with the same diffusion and $n_{W,p}$ for both cases.
 Convection needed to match radiation data near $p=0.6$ changes with atomic models: V⁸⁹, V^{TP}
 JETTO-NEO+BG8 [5-9] used to calculate theory-based transport coefficients, for comparison.



Due to flatness of n_e near $p=0.6$, the V⁸⁹ profile needed to reproduce the structure of P_{rad} in that region provides the best match to the theoretical pinch.

SUMMARY, PRELIMINARY CONCLUSIONS, FURTHER WORK

- W injection can produce detachment and a radiating mantle in the outer 10 cm inside the separatrix.
- Severely cooled edge allows deep penetration of neutrals, raising n_e & producing hollow n_e profiles.
- W motion along and across field lines has similar time scales: plumes seen in bolometry. This effect scales with machine size squared.
- Sawteeth can transport W inward to the core and back out.
- Work ongoing to improve ADAS predictions of W radiation in this relevant T_e range.
- Transport models required to match observed radiation data require less manipulation of the convection terms compared to theory when there is structure in $L_W(T_e)$ in the 0.3-1 keV range
- We hope to carry out more W ablation experiments into non-sawtooth plasmas

References:

[1] Beurskens Plasma Phys. Control. Fusion 55 (2013) t24043, [2] E. Stefanikova, 42nd EPS Conf. on Plasma Physics, Lisbon 2015, P2.131 [3] T. Pütterich et al. Plasma Phys. Control. Fusion 59 (2008) 085016 [4] T. Pütterich et al. Nuclear Fusion 50 (2010) 025012 [5] M. Erba et al. Nucl. Fusion 38 1013 (1998) [6] E. A. Belli, J. Candy PFCF 50 095010 (2008). [7] E. A. Belli, J. Candy PFCF 54 015015 (2012) [8] G. Cenacchi and A. Taroni, Report JET-IR 88 03 (1988). [9] G. Cenacchi and A. Taroni Rapporto ENEA RT/718 88 5 (1988).



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